

# TUTORIAL 05:

## NUMERICAL ASPECT II: CONSISTENCY AND STABILITY

In this tutorial, we will explore the consistency and stability of a numerical scheme. We will focus on a simple case study: the one-dimensional advection problem. Using Taylor expansion approximation, we will define the order of a numerical scheme and test its stability. We will reveal the famous CFL condition.

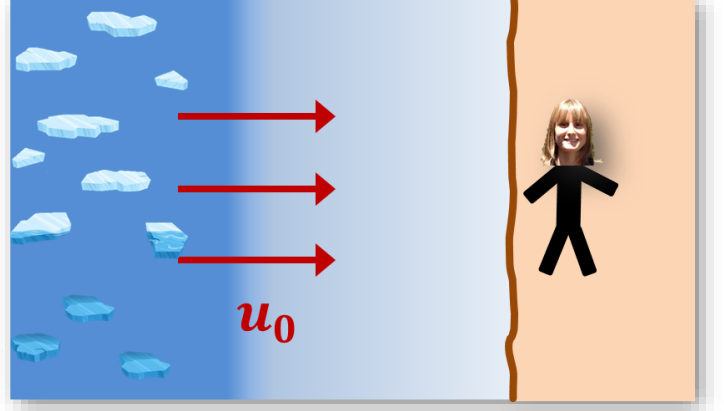
### 1: The 1D advection equation

→ From CROCO 3D temperature equation:

$$\frac{\partial T}{\partial t} + \mathbf{u} \nabla T = \nabla_h (K_{Th} \nabla_h T) + \frac{\partial}{\partial z} \left( K_{Tv} \frac{\partial T}{\partial z} \right)$$

→ We simplify the processes at work by studying a simple case study, where:

- there is no surface forcing (adiabatic).
- there is a constant current directed toward the shore  $u_0$  (homogeneous in  $y$ ).
- there is no variation of temperature with depth (barotropic case), i.e. we can cross-out the vertical turbulent diffusion term.
- there is no horizontal diffusion.



→ From the 3D temperature, we need to solve the 1D advection equation:

$$\frac{\partial T}{\partial t} + u_0 \frac{\partial T}{\partial x} = 0 \quad x \in [0, L], \quad t \in [0, T] \quad (1)$$

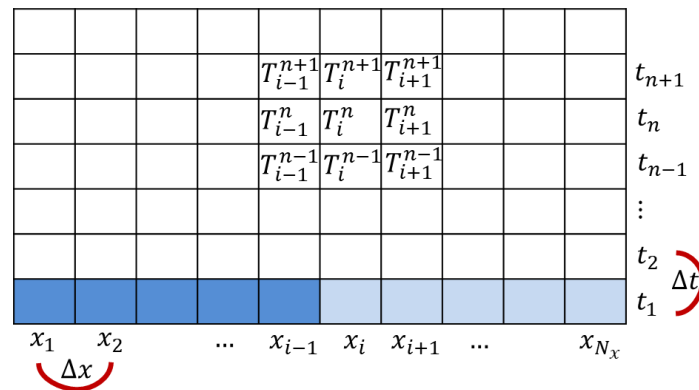
→ There are only a first-order derivatives in time and space.

→ The initial conditions that portray this temperature front are known. The constant parameter  $u_0$  (the current advecting the cold condition toward the coast) must be given.

### 2: Consistency of a numerical scheme

→ Same as in #TUTORIAL03, we work on a discretized model grid. We replace the continuous domain  $[0, L] \times [0, T]$  by a set of **equally spaced mesh points**, such that:

$$x_i = i\Delta x, i = 1, \dots, N_x \quad \text{and} \quad t_n = n\Delta t, n = 1, \dots, N_t$$



→ We need to find a **consistent** approximation for the equation derivatives:  $\frac{\partial T}{\partial t}$  and  $\frac{\partial T}{\partial x}$  on our model grid. This means that the error between the discretized and the real solution approaches 0.

→ In order to quantify the error we make by solving any equation on a spatial and temporal discretised grid, we use the Taylor series expansion of a continuous function  $f$  at a point  $x_0$  close to a reference point  $x$ :

$$f(x_0) = f(x) + \frac{f'(x)}{1!}(x_0 - x) + \frac{f''(x)}{2!}(x_0 - x)^2 + \dots + \frac{f^n(x)}{n!}(x_0 - x)^n + R(x)$$

↪ If  $x$  is close to  $x_0$ , such that  $x_0 = x + \Delta x$ , we can write:

$$f(x + \Delta x) = f(x) + \frac{f'(x)}{1!}\Delta x + \frac{f''(x)}{2!}\Delta x^2 + \dots + \frac{f^n(x)}{n!}\Delta x^n + R(x)$$

→ Let discretize  $\frac{\partial T}{\partial x}$ . There are 3 different numerical schemes:

❶ The **downstream** (Euler) scheme:  $\frac{\partial T}{\partial x} = \underline{\hspace{2cm}}$

❷ The **upstream** scheme:  $\frac{\partial T}{\partial x} = \underline{\hspace{2cm}}$

❸ The **centered** scheme:  $\frac{\partial T}{\partial x} = \underline{\hspace{2cm}}$



➤ Estimation of the error we make when we choose the downstream scheme (❶):

$$T(x + \Delta x) = T(x) + \frac{T'(x)}{1!}\Delta x + \frac{T''(x)}{2!}\Delta x^2 + \dots$$

$T'(x) = \underline{\hspace{2cm}}$

➤ Estimation of the error we make when we choose the upstream scheme (2):

$$T(x - \Delta x) = T(x) - \frac{T'(x)}{1!} \Delta x + \frac{T''(x)}{2!} \Delta x^2 + \dots$$

$$T'(x) = \underline{\hspace{5cm}}$$

➤ Estimation of the error we make when we choose the centered scheme (3):

$$T'(x) = \underline{\hspace{5cm}}$$

↪ With the centered scheme, the first-order derivative is better resolved than with the first order schemes.

⇒ The centered scheme is better than upstream and downstream schemes, because the **truncation error** is smaller. To improve it, you can increase your resolution ( $\Delta x \searrow$ ) or use higher-order schemes.

### 3: Stability and convergence of a numerical scheme

→ The most important characteristics of a numerical scheme are:

- Its **consistency**, i.e. consistent approximation for the derivative in the equations (truncation error  $\searrow 0$ ). This is a condition in space.

- Its **stability**, i.e. does the error amplify in time? We do not want that the error increase with time. If this the case, there will be a numerical explosion (a blow-up), and the model will stop.

⇒ If both conditions are respected (consistency and stability) then the discrete solution **converges** toward the real solution.

❶ We will test the stability of a **downstream** scheme for both:  $\frac{\partial T}{\partial t}$  and  $\frac{\partial T}{\partial x}$ , such that:

$$\frac{\partial T}{\partial t} \approx \frac{T(t + \Delta t) - T(t)}{\Delta t} = \underline{\hspace{4cm}}$$

$$\frac{\partial T}{\partial x} \approx \frac{T(x + \Delta x) - T(x)}{\Delta x} = \underline{\hspace{4cm}}$$

→ We inject this formulation into the 1D-advection equation. This leads to:

$$\frac{\partial T}{\partial t} + u_0 \frac{\partial T}{\partial x} = 0 \quad \rightarrow$$

→

↪ This gives  $T$  at time  $t + \Delta t$  as a function of  $T$  at time  $t$ . This is an **explicit method**. It is easy to solve

→ We will perform a **von Neumann** stability analysis of our explicit solution.

↪ For this we use wave-like structure for  $T(x)$  using complex form:  $T_n = \hat{T}_n e^{ikx}$

- $e^{ikx}$  is a wavy pattern that repeats indefinitely ( $k$  provide information about its zonal extension).

- $\hat{T}_n$  is :

→ We rewrite our explicit solution using this new notation.

$$\hat{T}_{n+1} e^{ikx} =$$

With  $C = \frac{u_0 \Delta t}{\Delta x} > 0$ , the Courant number.

→ We now define the amplification A, such that:

$$A = \frac{\hat{T}_{n+1}}{\hat{T}_n}$$

↪ We want  $A < 1$ , because we do not want the amplitude of oscillation to increase over time, otherwise the solution would explode.

$$A = \frac{\hat{T}_{n+1}}{\hat{T}_n} =$$

$$\|A\|^2 =$$

⇒  $\|A\| > 1$ . This means that the solution increases over time. This scheme is unstable. The downstream scheme is not a good choice. I will never know if I can go to the beach tomorrow

② We will use the downstream scheme in space, and the upstream scheme in time. This is the upwind scheme:

$$\frac{\partial T}{\partial t} \approx \frac{T(t + \Delta t) - T(t)}{\Delta t} = \underline{\hspace{2cm}}$$

$$\frac{\partial T}{\partial x} \approx \frac{T(x) - T(x - \Delta x)}{\Delta x} = \underline{\hspace{2cm}}$$

↪ This leads to:

↪ We adopt the complex form:  $T_n = \widehat{T}_n e^{ik(x)}$ . We obtain:

↪ We again define the amplification  $A = \frac{\widehat{T}_{n+1}}{\widehat{T}_n}$ , such that:

$$A = \frac{\widehat{T}_{n+1}}{\widehat{T}_n} =$$

$$\|A\|^2 =$$

⇒ This upwind scheme is conditionally stable. Is is stable if  $C = \frac{u_0 \Delta t}{\Delta x} < 1$ . This is the famous CFL condition.

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=  STEP2D:  ABNORMAL JOB END
=          BLOW UP
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