

TUTORIAL 05:

NUMERICAL ASPECT II: CONSISTENCY AND STABILITY

In this tutorial, we will explore the consistency and stability of a numerical scheme. We will focus on a simple case study: the one-dimensional advection problem. Using Taylor expansion approximation, we will define the order of a numerical scheme and test its stability. We will reveal the famous CFL condition.

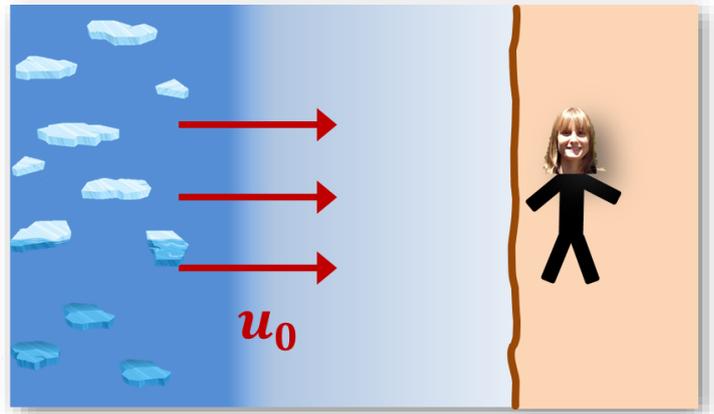
1: The 1D advection equation

→ From CROCO 3D temperature equation:

$$\frac{\partial T}{\partial t} + \mathbf{u}\nabla T = \nabla_h(K_{Th}\nabla_h T) + \frac{\partial}{\partial z}\left(K_{Tv}\frac{\partial T}{\partial z}\right)$$

↳ We simplify the processes at work by studying a simple case study, where:

- there is no surface forcing (adiabatic).
- there is a constant current directed toward the shore u_0 (homogeneous in y).
- there is no variation of temperature with depth (barotropic case), i.e. we can cross-out the vertical turbulent diffusion term.
- there is no horizontal diffusion.



→ From the 3D temperature, we need to solve the 1D advection equation:

$$\frac{\partial T}{\partial t} + u_0 \frac{\partial T}{\partial x} = 0 \quad x \in [0, L], \quad t \in [0, T] \quad (1)$$

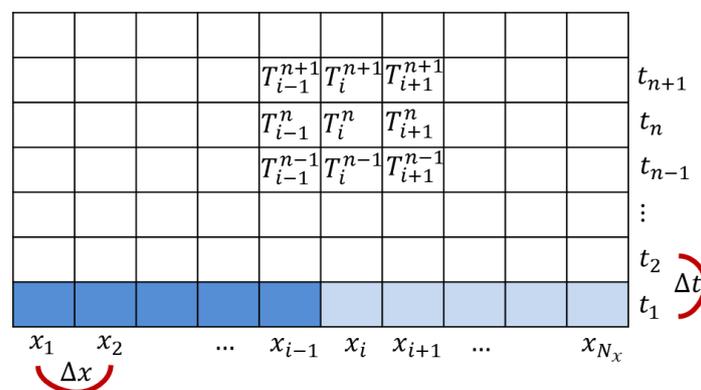
↳ There are only first-order derivatives in time and space.

↳ The initial conditions that portray this temperature front are known. The constant parameter u_0 (the current adveting the cold condition toward the coast) must be given.

2: Consistency of a numerical scheme

→ Same as in #TUTORIAL03, we work on a discretized model grid. We replace the continuous domain $[0, L] \times [0, T]$ by a set of **equally spaced mesh points**, such that:

$$x_i = i\Delta x, i = 1, \dots, N_x \quad \text{and} \quad t_n = n\Delta t, n = 1, \dots, N_t$$



→ We need to find a **consistent** approximation for the equation derivatives: $\frac{\partial T}{\partial t}$ and $\frac{\partial T}{\partial x}$ on our model grid. This means that the error between the discretized and the real solution approaches 0.

→ In order to quantify the error we make by solving any equation on a spatial and temporal discretised grid, we use the Taylor series expansion of a continuous function f at a point x_0 close to a reference point x :

$$f(x_0) = f(x) + \frac{f'(x)}{1!}(x_0 - x) + \frac{f''(x)}{2!}(x_0 - x)^2 + \dots + \frac{f^n(x)}{n!}(x_0 - x)^n + R(x)$$

↪ If x is close to x_0 , such that $x_0 = x + \Delta x$, we can write:

$$f(x + \Delta x) = f(x) + \frac{f'(x)}{1!}\Delta x + \frac{f''(x)}{2!}\Delta x^2 + \dots + \frac{f^n(x)}{n!}\Delta x^n + R(x)$$

→ Let discretize $\frac{\partial T}{\partial x}$. There are 3 different numerical schemes:

① The **downstream** (Euler) scheme: $\frac{\partial T}{\partial x} =$ _____

② The **upstream** scheme: $\frac{\partial T}{\partial x} =$ _____

③ The **centered** scheme: $\frac{\partial T}{\partial x} =$ _____



➤ Estimation of the error we make when we choose the downstream scheme (①):

$$T(x + \Delta x) = T(x) + \frac{T'(x)}{1!}\Delta x + \frac{T''(x)}{2!}\Delta x^2 + \dots$$

$T'(x) =$ _____

➤ Estimation of the error we make when we choose the upstream scheme (2):

$$T(x - \Delta x) = T(x) - \frac{T'(x)}{1!} \Delta x + \frac{T''(x)}{2!} \Delta x^2 + \dots$$

$T'(x) =$ _____

➤ Estimation of the error we make when we choose the centered scheme (3):

$T'(x) =$ _____

↳ With the centered scheme, the first-order derivative is better resolved than with the first order schemes.

⇒ The centered scheme is better than upstream and downstream schemes, because the **truncation error** is smaller. To improve it, you can increase your resolution ($\Delta x \searrow$) or use higher-order schemes.

3: Stability and convergence of a numerical scheme

→ The most important characteristics of a numerical scheme are:

- Its **consistency**, i.e. consistent approximation for the derivative in the equations (truncation error $\searrow 0$). This is a condition in space.

- Its **stability**, i.e. does the error amplify in time? We do not want that the error increase with time. If this the case, there will be a numerical explosion (a blow-up), and the model will stop.

⇒ If both conditions are respected (consistency and stability) then the discrete solution **converges** toward the real solution.

❶ We will test the stability of a **downstream** scheme for both: $\frac{\partial T}{\partial t}$ and $\frac{\partial T}{\partial x}$, such that:

$$\frac{\partial T}{\partial t} \approx \frac{T(t + \Delta t) - T(t)}{\Delta t} = \underline{\hspace{10em}}$$

$$\frac{\partial T}{\partial x} \approx \frac{T(x + \Delta x) - T(x)}{\Delta x} = \underline{\hspace{10em}}$$

→ We inject this formulation into the 1D-advection equation. This leads to:

$$\frac{\partial T}{\partial t} + u_0 \frac{\partial T}{\partial x} = 0 \quad \rightarrow$$

→

↪ This gives T at time $t + \Delta t$ as a function of T at time t . This is an **explicit method**. It is easy to solve

→ We will perform a **von Neumann** stability analysis of our explicit solution.

↪ For this we use wave-like structure for $T(x)$ using complex form: $T_n = \hat{T}_n e^{ikx}$

- e^{ikx} is a wavy pattern that repeats indefinitely (k provide information about its zonal extension).

- \hat{T}_n is :

→ We rewrite our explicit solution using this new notation.

$$\hat{T}_{n+1} e^{ikx} =$$

With $C = \frac{u_0 \Delta t}{\Delta x} > 0$, the Courant number.

→ We now define the amplification A, such that:

$$A = \frac{\hat{T}_{n+1}}{\hat{T}_n}$$

↳ We want $A < 1$, because we do not want the amplitude of oscillation to increase over time, otherwise the solution would explode.

$$A = \frac{\hat{T}_{n+1}}{\hat{T}_n} =$$

$$\|A\|^2 =$$

⇒ $\|A\| > 1$. This means that the solution increases over time. This scheme is unstable. The downstream scheme is not a good choice. I will never know if I can go to the beach tomorrow

② We will use the downstream scheme in space, and the upstream scheme in time. This is the upwind scheme:

$$\frac{\partial T}{\partial t} \approx \frac{T(t + \Delta t) - T(t)}{\Delta t} = \underline{\hspace{10em}}$$

$$\frac{\partial T}{\partial x} \approx \frac{T(x) - T(x - \Delta x)}{\Delta x} = \underline{\hspace{10em}}$$

↳ This leads to:

↳ We adopt the complex form: $T_n = \widehat{T}_n e^{ik(x)}$. We obtain:

↳ We again define the amplification $A = \frac{\widehat{T}_{n+1}}{\widehat{T}_n}$, such that:

$$A = \frac{\widehat{T}_{n+1}}{\widehat{T}_n} =$$

$$\|A\|^2 =$$

⇒ This upwind scheme is conditionally stable. It is stable if $C = \frac{u_0 \Delta t}{\Delta x} < 1$. This is the famous CFL condition.

